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# Gravitational Wave Detectors

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Postgrad Lectures 2013

# Outline

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- Interferometer basics
  - Recap of GW basics
  - Detecting gravitational waves
  - Interferometers
  - GW detectors
- Noise sources
  - Seismic noise
  - Optimal readout noise
  - Thermal noise
- Transient noise

# Who am I?

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- Thomas Adams
- 4<sup>th</sup> year PhD student working with Patrick
  - Searches for gravitational wave bursts
  - Characterisation of the GEO 600 detector
- Spent half my PhD on site at the British - German GEO 600 detector near Hannover

# Recap of gravitational waves

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- Distance between two points:

$$ds^2 = \left( dx_1^2 + dx_2^2 + dx_3^2 \right)$$

- Can rewrite this in General Relativity as

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$

- Where  $g$  is the metric (identity for flat space) and  $dx$  is the displacement vector
- Includes new term for displacement in time

# Recap of gravitational waves

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- Gravitational waves are perturbation of the space-time metric

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

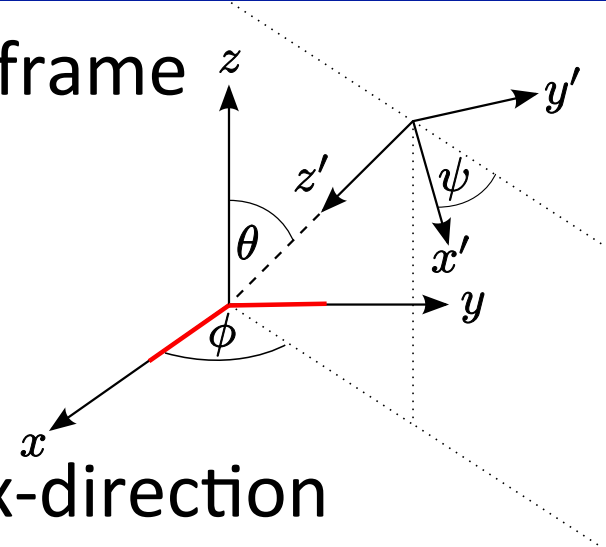
- In the simplest form, waves travelling in z-direction take this form:

$$h_{\mu\nu} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & h_+ & h_\times & 0 \\ 0 & h_\times & -h_+ & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

# Detecting gravitational waves

- Rotation gives simple form in lab frame

$$h_{\mu\nu} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & h_{xx} & h_{xy} & h_{xz} \\ 0 & h_{yx} & h_{yy} & h_{yz} \\ 0 & h_{zx} & h_{zy} & h_{zz} \end{pmatrix}$$



- Consider light beam travelling in x-direction (down the lab) between two test masses, and GW travelling in z-direction (vertically down). Distance between them changes:

$$\begin{aligned} ds^2 = g_{\mu\nu} dx^\mu dx^\nu &= (\eta^{\mu\nu} + h_{\mu\nu}) dx^\mu dx^\nu \\ &= -c^2 dt^2 + dx^2 (1 + h_{xx}^{TT}) \end{aligned}$$

# Detecting gravitational waves

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- Consider the following experiment
  - Shine a laser down the lab (in the x-direction)
  - Time how long it takes to hit the test mass

- For a laser beam:

$$ds^2 = 0$$

- so the distance to the end of the lab changes

$$L_x = \int_0^L c dt = \int_0^L \sqrt{1 + h_{xx}} \simeq \left(1 + \frac{1}{2}h_{xx}\right) L$$

- We can measure the light travel time change to detect and study gravitational waves directly

# Detecting gravitational waves

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- If we do this in both the x- and y-directions we can measure gravitational wave strain

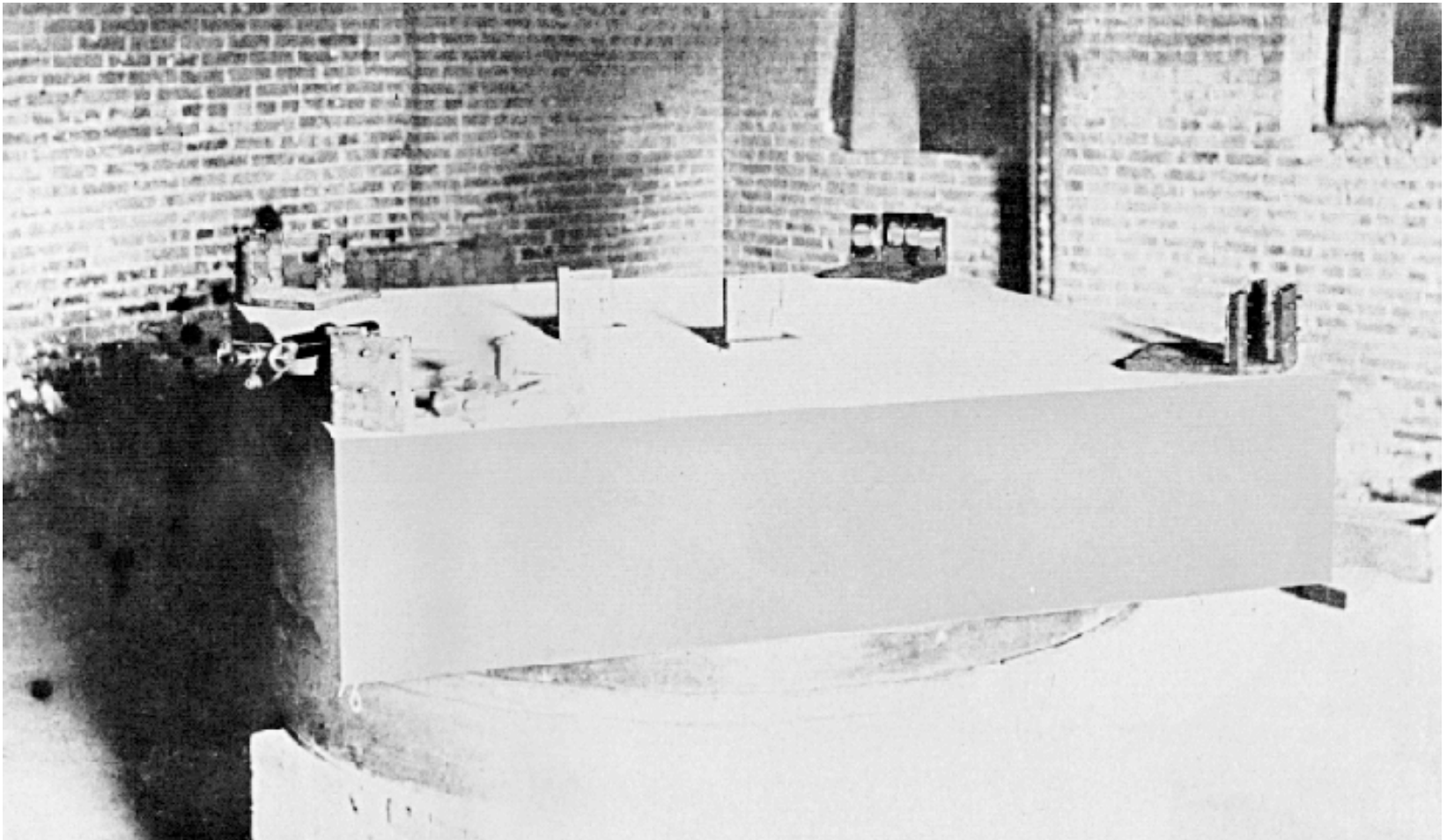
$$h(t) = \frac{\Delta L}{L} = \frac{L_x - L_y}{L} = \frac{1}{2} (h_{xx} - h_{yy})$$

- Conveniently, such experiments have been around for more than a hundred years



# Interferometers

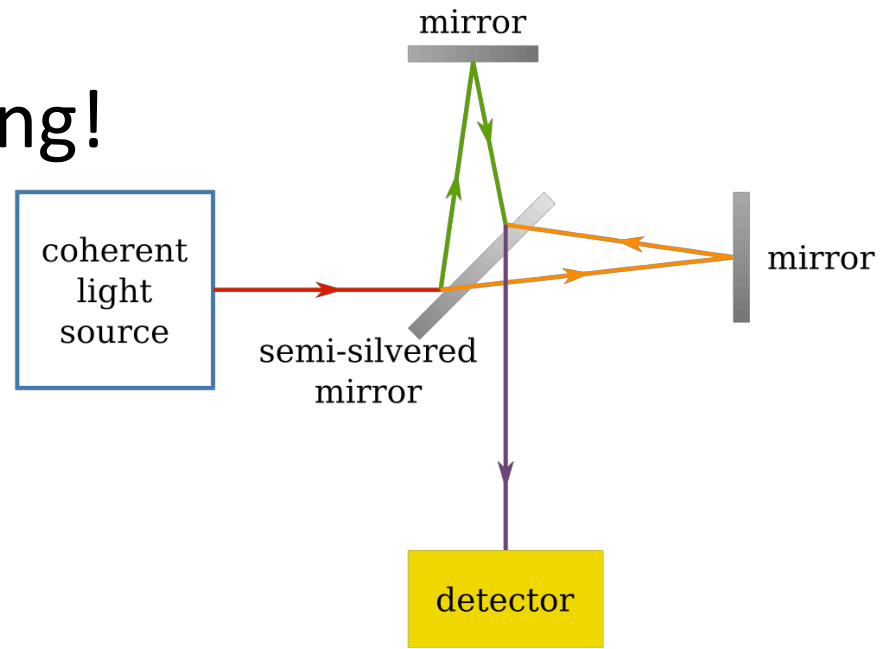
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# Interferometers

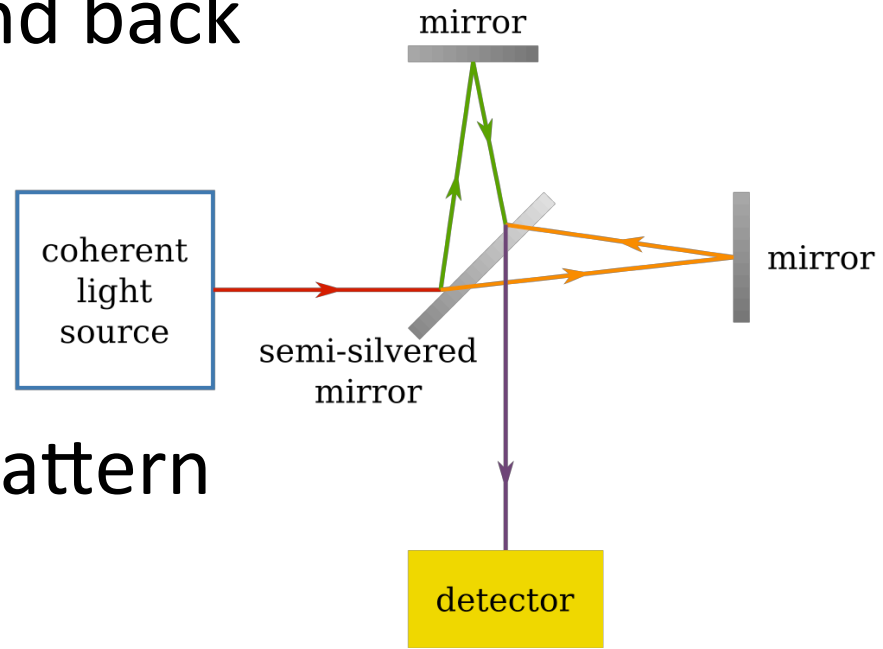
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- Michelson-Morley designed this experiment to detect the  $\text{\AA}$ ether in 1887
- Proved themselves wrong!



# Interferometers

- Shine laser light onto partially reflecting mirror – “beam splitter”
- Send down two arms and back
- Recombine at beam splitter and exit at output photo detector
- Measure interference pattern to accuracy of  $\lambda/20$



# GW detectors



American LIGO at Hanford



American LIGO at Livingston



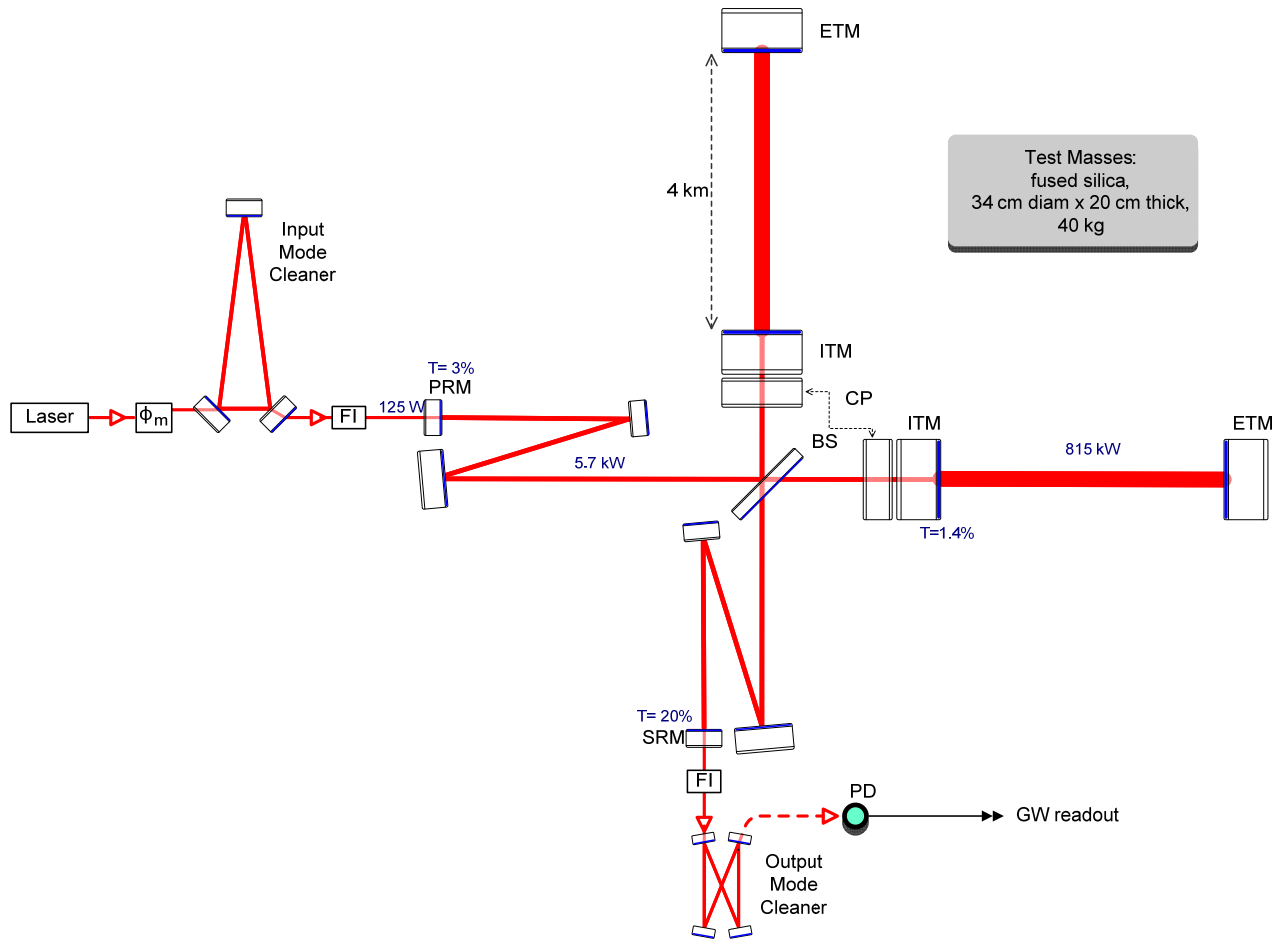
British-German GEO

GEO 600



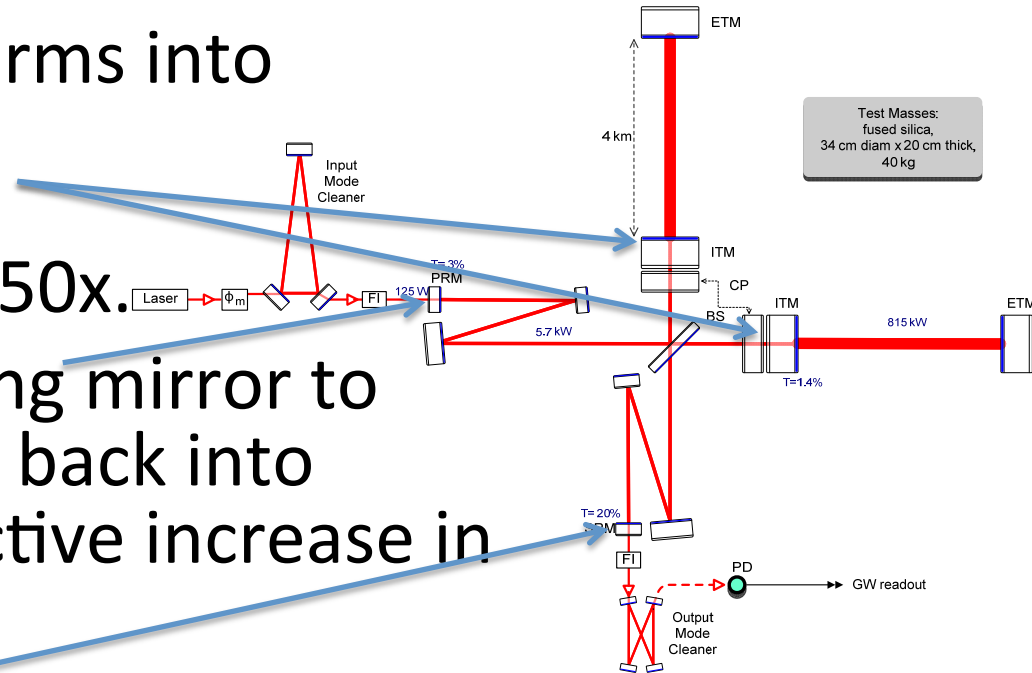
French-Italian Virgo near PISA

# GW detectors



# GW detectors

- Added mirrors turn arms into Fabry-Perot cavities. Effective increase in arm length of 30x-50x.
- Added power recycling mirror to return reflected light back into interferometer. Effective increase in power by 20x.
- Added signal recycling mirror to return output light (signal) into interferometer. Effective increase in signal time in interferometer.



# GW detectors

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- Have global network of large scale detectors
- Laser Interferometer Gravitational-Wave Observatory (LIGO):
  - Two detectors at Hanford, WA
    - ‘H1’ – 4 kilometer Fabry-Perot arm cavities
    - ‘H2’ – 2 km FP cavities
  - One detector at Livingston, LA
    - ‘L1’ – 4 km FP cavities
- Virgo at Pisa, Italy, with 3 km FP cavities
- GEO600 at Hannover, Germany with 600m folded arms
  - No FP cavities, arms fold back on themselves



# GW detector sensitivity

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- Optimum sensitivity is achieved when light takes  $\frac{1}{2}$  period of GW to traverse the arm
- For a GW at 200Hz, this corresponds to arm length of 500km...
- Fabry-Perot cavities store light in arms for  $\sim 50$  round trips, LIGO with FP has effective arm length 200km.



# GW detector sensitivity

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- Consider a GW travelling in z-direction (i.e. landing on top of detector).
- We have  $h_{xx} = -h_{yy} = h$ , and recall for a single trip down the x-arm

$$c\Delta t = \left(1 + \frac{1}{2}h_{xx}\right) L$$

- So, for round trip, time delay between signals in two arms is:

$$\Delta t = \frac{2}{c} \left( \left(1 + \frac{1}{2}h_{xx}\right) L - \left(1 - \frac{1}{2}h_{yy}\right) L \right)$$

# GW detector sensitivity

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- Solving that gives a time-delay

$$\Delta t = \frac{2h(t)L}{c}$$

or a phase shift of

$$\Delta\phi = \frac{4\pi h(t)L}{\lambda}$$

- If we plug in realistic numbers,  $\lambda = 1064\text{nm}$ ,  $L = 4\text{km}$ ,  $h = 10^{-21}$  we get:

$$\Delta L = 2 \times 10^{-18}\text{m}, \quad \Delta\phi = 2 \times 10^{-9}$$

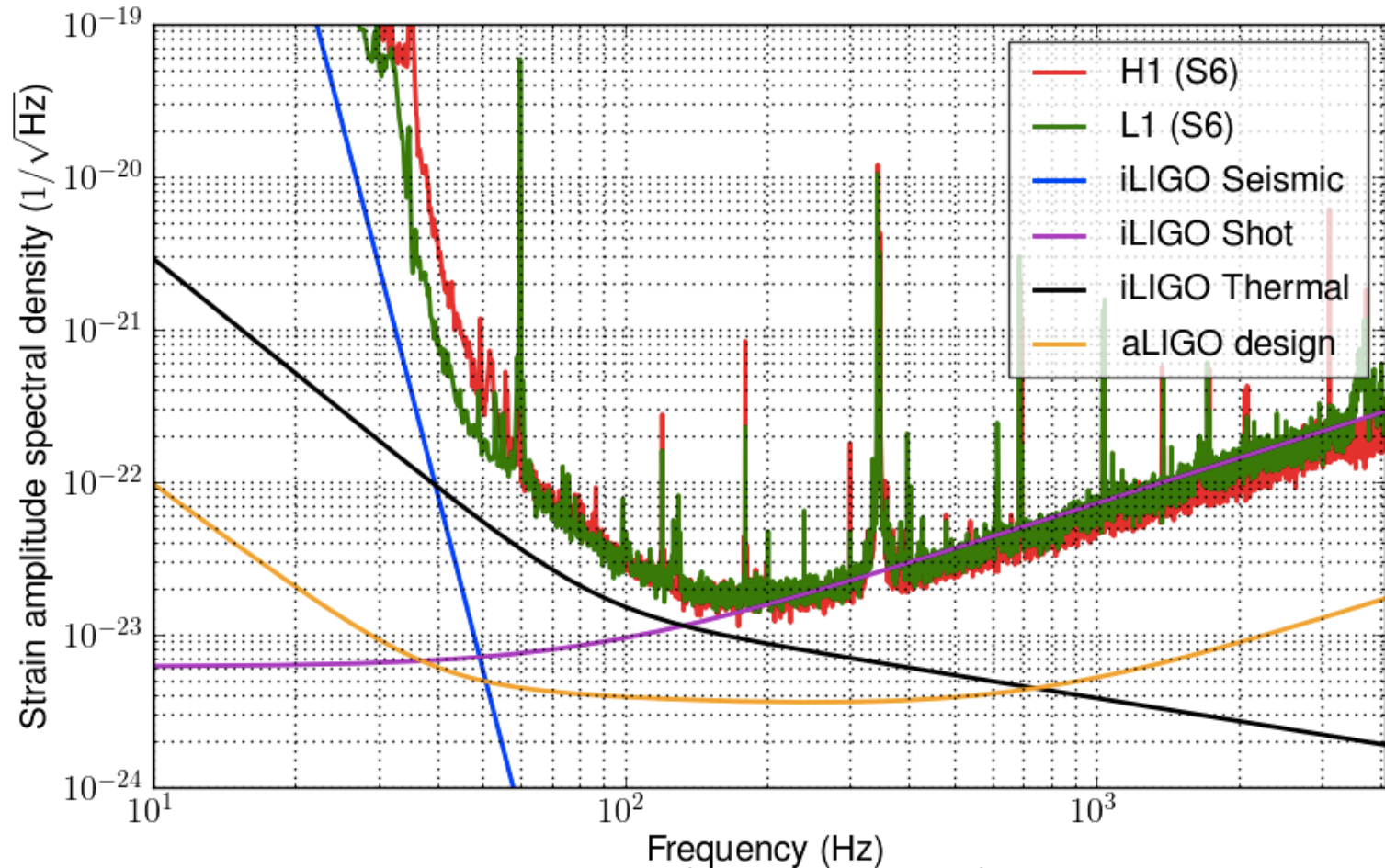
# GW detector sensitivity

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- Patrick Sutton's reality check:
  - Sensitivity  $\Delta L = 2 \times 10^{-18}\text{m}$   
 $\Delta\varphi = 2 \times 10^{-9}$
- These are:
  - $10^{10}$  times smaller than 1 MM fringe
  - $10^{12}$  times smaller than laser wavelength
  - $10^9$  times smaller than atoms in mirrors whose position we try to measure
  - $10^{12}$  times smaller than seismic motion shaking mirrors



# GW detector sensitivity



- It works! ...but it cost \$250M + \$250M for aLIGO

# Seismic noise

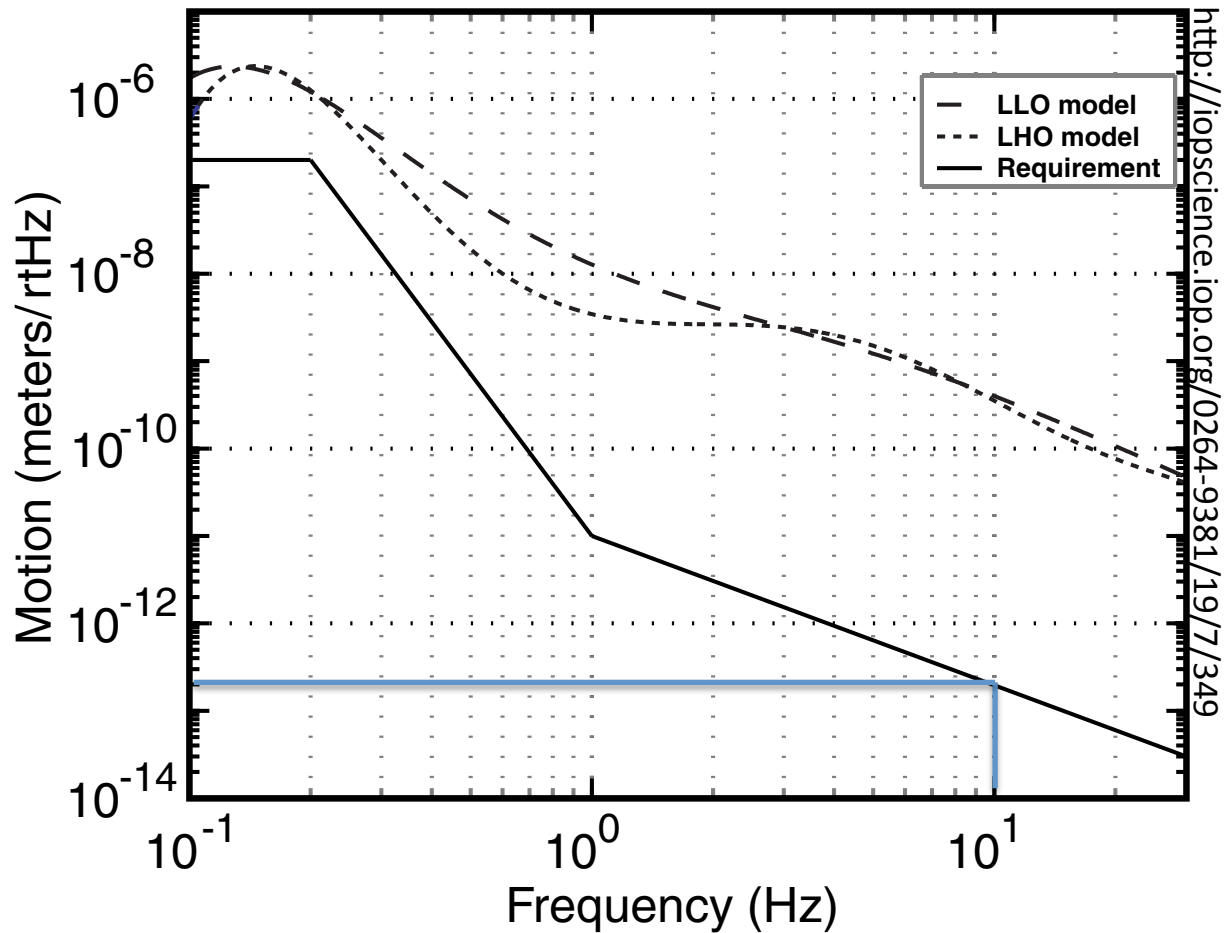
- Shaking of ground due to earthquakes, weather, human activity

Frequency (Hz)	Distance (km)	Source
0.01 – 1	$10^3$	Distant earthquakes Microseism
1 – 3	$10^1$	Far anthropogenic noise Close earthquakes Wind
3 – 10	$10^0$	Anthropogenic noise Wind
10 – 30	$10^{-1}$	Close anthropogenic noise

<http://arxiv.org/abs/1108.0312>

**Table 1.** Description of the main seismic frequency bands and their sources

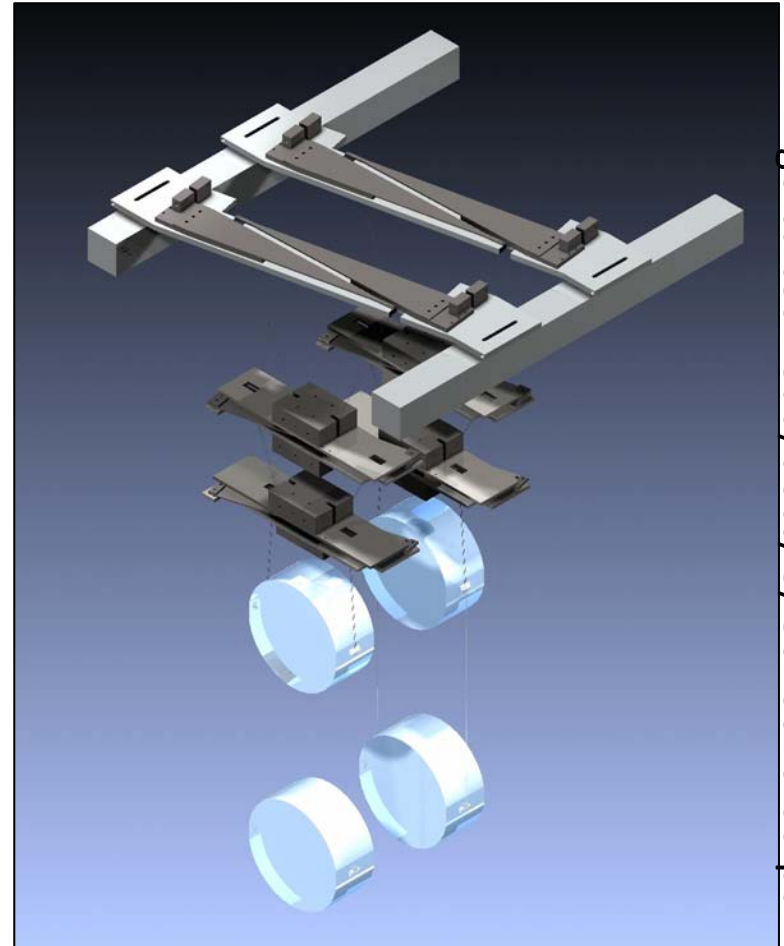
# Seismic noise



At 10 Hz, seismic motion is several orders of magnitude above GW spectrum

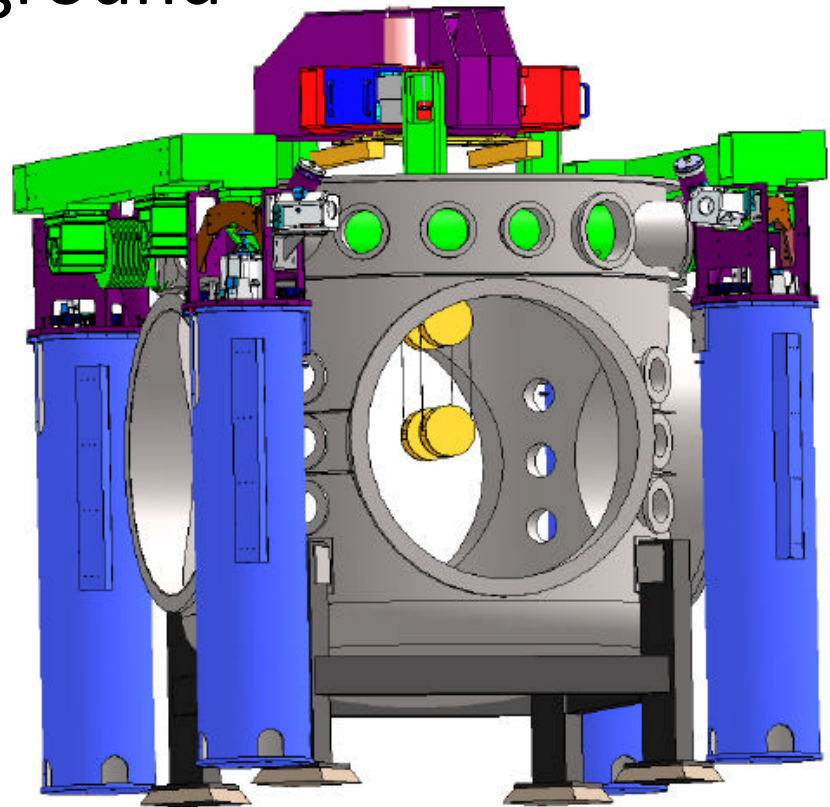
# Passive seismic isolation

- Need to suppress motion by factor of  $10^8$
- Mirrors suspended as quadruple pendula
- 40 kg silica test masses
- Lab floor on separate slab



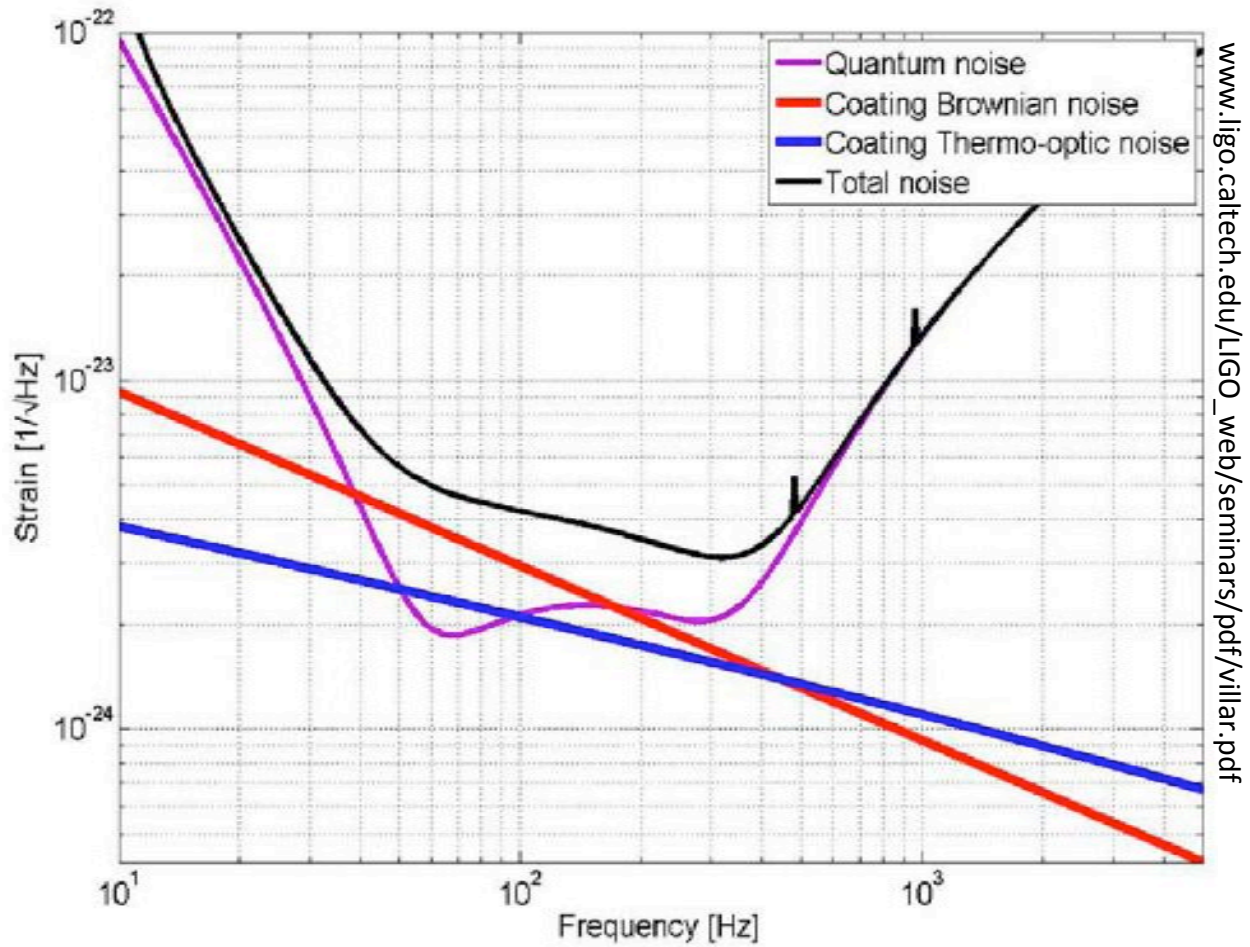
# Active seismic isolation

- Use sensors to detect ground motion and correct accordingly
- Hydraulic actuators reduce factor of 10 at low frequency ( $<2$  Hz)
- EM actuators give factor of 30 at 10Hz





# Thermal noise



www.ligo.caltech.edu/LIGO\_web/seminars/pdf/villar.pdf

# Thermal noise

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- All detector parts are subject to thermal noise:
  - Brownian motion (random motion of atoms in the mirrors)
  - Vibrational modes (resonances) of the suspension wires
  - Dissipation from friction in the wires
- Recall equipartition theorem: each degree of freedom contributes energy of  $\frac{1}{2}k_B T$ . For 40 kg mass suspended from 1m wire:

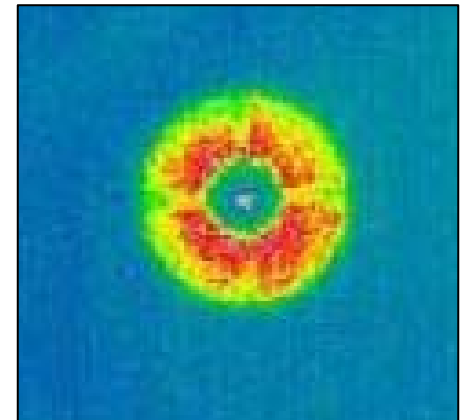
$$\Delta L_{\text{rms}} = \sqrt{\frac{k_B T l}{m g}} \simeq 3 \times 10^{-12} \text{ m}$$

- $10^7$  times larger than motion due to GW

# Thermal noise

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- Further thermal noise due to heat transfer from laser to mirrors
- Mitigate thermal noise with:
  - Fused silica wires – less friction, less noise
  - Better connection to mirror
  - Better coatings for mirrors
  - Heavier mirrors
  - Thermal compensation



# Quantum noise – shot noise

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- Quantum noise is introduced by the fluctuating number of photons arriving at the output port
- The uncertainty in the phase of a laser beam due to quantization of light into photons is called shot noise.
- Dominant noise source at high frequency
- If error in photon number is  $\sigma_N$ , and error in phase is  $\sigma_\phi$  HUP:

$$\sigma_N \sigma_\phi = 1$$

# Quantum noise – shot noise

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- Number of photons per second:

$$N_\gamma = \frac{P}{E} = \frac{P \lambda}{hc} = 50 \frac{200\text{W } 1064\text{nm}}{hc} \simeq 6 \times 10^{22} \text{ s}^{-1}$$

- So error (assuming Poisson statistics):

$$\sigma_N = \sqrt{N} = 2.5 \times 10^{11}$$

- And so the phase error is:

$$\sigma_\phi = \frac{1}{\sigma_N} = 4 \times 10^{-12} \text{ rad}$$

- This is equivalent to an arm length change of:

$$\Delta L = \frac{1064 \text{ nm}}{2\pi} 4 \times 10^{-12} = 6 \times 10^{-19} \text{ m}$$

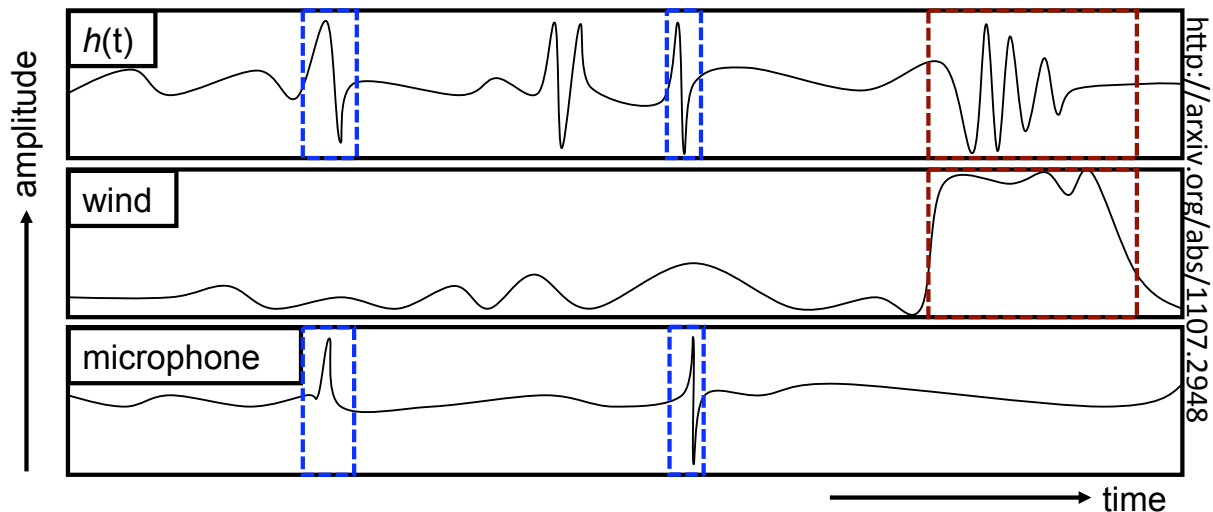
# Quantum noise – radiation pressure

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- Quanta in laser beam carry energy and momentum, which pass to mirrors upon transmission/reflection
- Higher power laser gives better sensitivity in areas (less shot noise), but higher radiation pressure noise
- Limiting noise source in 10-50 Hz band.

# Transient noise

- Limiting sources are stationary (time and freq)
  - form baseline for sensitivity
- Many searches for short-duration GW events
- Also short-duration noise events that mask/mimic GWs



# Summary

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- Differential arm motion can detect GW!
- Large-scale interferometers have been built for this purpose
  - Can detect differential motion to  $10^{-18}$  m!
- Subject to noise sources
  - Seismic noise due to ground motion (low frequency)
  - Thermal noise in suspension and mirrors (mid frequency)
  - Radiation pressure noise (mid) and shot noise (high) due to laser
- Subject to transient events that can mimic GW